

FUNCTION DYNAMICS

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Abstract. We show mathematical structure of the function dynamics, i.e., the dynamics of interval maps $f_{n+1} = (1 - \varepsilon)f_n + \varepsilon f_n \circ f_n$ and clarify the types of fixed points, the self-referential structure and the hierarchical structure.

1. Introduction. Since about thirty years ago the late Professor Yamaguti had continued to recommend young researchers to find and study “new mathematics in phenomena”. Figure 1.1 below shows the phenomena we study in the present paper.

Let I be an interval and for a given map $f : I \rightarrow I$ let us define a new map $\Phi_\varepsilon(f) : I \rightarrow I$ by

$$(1.1) \quad \Phi_\varepsilon(f) = (1 - \varepsilon)f + \varepsilon f \circ f,$$

where $0 < \varepsilon \leq 1$. Given $f_0 : I \rightarrow I$ we consider the function dynamics defined by

$$(1.2) \quad f_{n+1} = \Phi_\varepsilon(f_n) \quad (n = 0, 1, 2, \dots).$$

The original motivation to study (1.1) of [1] is mentioned in Section 5, but the motivation of the present paper consists in the study of (1.1) as an infinite dimensional dynamical system. For $\varepsilon = 1$, $f_{n+1} = f_n \circ f_n$ is nothing but 2^n -th iteration of the map f_1 . Therefore, one might expect very much complicated and chaotic behaviors in (1.1). However, the simulations in [1,2] for $\varepsilon < 1$ indicate that (1.1) can exhibit rather simple behaviors with hierarchical and self-referential structures, which we will prove in a rigorous manner in the present paper.

The above dynamics can also be written as

$$(1.3) \quad f_{n+1} = g_n \circ f_n$$

where g_n is defined from f_n by $g_n(x) = (1 - \varepsilon)x + \varepsilon f_n(x)$. The structure that f_n gives g_n and f_n is evolved by g_n is a key to the emergence of what we call self-referential structure.

Figure 1.1 shows two typical examples of the phenomena observed in function dynamics as $n \rightarrow \infty$ with $f_0(x) = rx(1 - x)$. In the simulation, we take a finite mesh size to compute the function $f_n(x)$, although the “phenomena” we discuss is not an artifact of the finite mesh, but they remain as mesh points are increased (or one can say that a piece-wise step function to approximate $f_0(x)$ with a small mesh size). As n goes large in the simulation the flat parts of the graphs grow up rapidly and they fill the whole interval in Figure 1.1(a) (within 100 simulation steps when mesh number = 4096). At each flat part $f_n(x)$ starts to be fixed in time within some time steps. There appear finer flat parts with smaller intervals when the initial r is larger or ε is

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smaller. Furthermore, there appear some complicated structure and some parts with irregular oscillation in time, near the end points of flat parts in Figure 1.1(b). Those phenomena as well as other structures and dynamics were reported and heuristically analyzed in [1] and [2].

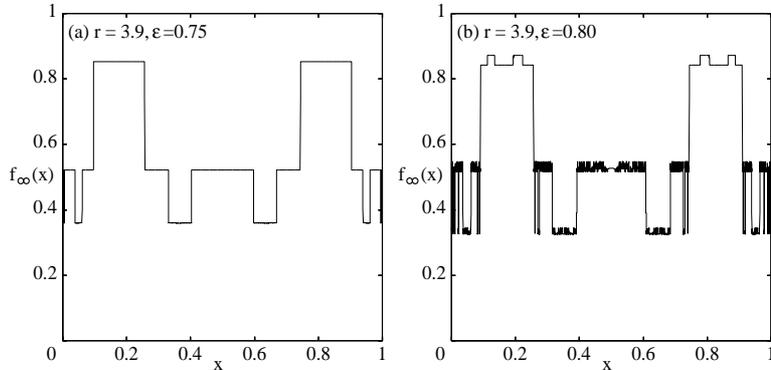


FIG. 1.1. Two typical results $f_\infty(x)$ of numerical simulation of function dynamics for $f_0(x) = rx(1-x)$. (a) $r = 3.9$, $\varepsilon = 0.75$. (b) $r = 3.9$, $\varepsilon = 0.8$. $f_n(x)$ with sufficiently large n is plotted. The simulation is carried out by using the mesh number 4096, i.e., by taking a piecewise step function approximating $f_0(x)$ with the mesh size $1/4096$.

In the present paper, we study the limit $f_\infty(x)$. In Section 2, the flat parts are explained. In Section 3 we study how the “self-reference” is represented within a hierarchical structure of the function $f_\infty(x)$.

In particular, there can exist trajectories such that

$$f_{n+1}(x) = g_\infty(f_n(x))$$

where g_∞ is the “generated map” in the terminology of [1], and some part of function is driven by other parts. In Section 4, we give a further example which shows “entangled hierarchy”, where the hierarchy of rules change dynamically in time. In the last section 5, the original motivation of the model are discussed, and the results are interpreted.

To close the introduction, we state some terminology. We denote the set of fixed points of f as $Fix(f) = \{x \in I \mid f(x) = x\}$. Take a fixed point $q \in Fix(f)$, and we call q a stable fixed point if there exists an open neighborhood U of q such that $U \supset f(U) \supset \dots \supset f^n(U)$ and $\bigcap_{n \geq 0} f^n(U) = \{q\}$. A semi-stable fixed point q is defined in a similar manner but U has the form $[q, q + \delta)$ or $(q - \delta, q]$. $Fix^{(s)}(f)$ denotes the set of stable or semi-stable fixed points. The basin of attraction $B(Q)$ of $Q \subset Fix^{(s)}(f)$ is defined as $\{x \in I \mid \lim_{n \rightarrow \infty} f^n(x) \in Q\}$.

2. The fixed point f_∞ . Our starting point of the study is to focus on those points x in the interval I where the limits

$$(2.1) \quad f_\infty(x) = \lim_{n \rightarrow \infty} f_n(x)$$

exist.

Optimists will take the formal limit of $f_{n+1} = (1 - \varepsilon)f_n + \varepsilon f_n \circ f_n$ to find the following relation independent of $\varepsilon > 0$:

$$(2.2) \quad f_\infty(x) = f_\infty(f_\infty(x)).$$

In other words, it is expected that the limit $f_\infty(x)$ is, if any, a fixed point of the map f_∞ and that f_∞ is a step function taking fixed point as its values. Simulations support this. Indeed, what we called the flat parts in Figure 1.1 form step functions. Note, however, that domain of f_∞ may not be the whole interval I .

The following is the mathematical statement for the above observation.

THEOREM 2.1. *For a given continuous map $f_0 : I \mapsto I$, there exist a non-empty subset Ω of the interval I and a map $f_\infty : \Omega \mapsto \Omega$ which satisfy the following properties:*

- (i) *For each $x \in \Omega$ the limit (2.1) exists and (2.2) holds.*
- (ii) *$\text{Fix}(f_\infty)$ is non-empty.*
- (iii) *$f_\infty(\Omega) \subset \text{Fix}(f_\infty)$. In other words, f_∞ is a step function on Ω outside a (possibly empty) subset of I where $f(x) = x$. More precisely, let $\Omega_q := f_\infty^{-1}(q) = \{x \in \Omega | f_\infty(x) = q\}$, we have a partition:*

$$\Omega = \bigcup_{q \in \text{Fix}(f_\infty)} \Omega_q.$$

REMARK 2.2. *We should mention here that we do not exclude the case when the map f_∞ is the identity map if it is restricted to a subinterval.*

As the proof below shows, we can take

$$(2.3) \quad \Omega = \bigcup_{n \geq 0} f_n^{-1}(\text{Fix}(f_n)).$$

In particular, $\Omega \neq \emptyset$ since $\text{Fix}(f_0) \neq \emptyset$ by the intermediate value theorem. Moreover, the set

$$(2.4) \quad \Omega^I = \bigcup_{n \geq 0} \text{Fix}(f_n)$$

coincides with $\text{Fix}(f_\infty)$. In [1], the point in Ω^I is called the fixed point of type-I and the point in $\Omega^{II} := \Omega \setminus \Omega^I$ is called the fixed point of type-II.

The set Ω_q is an at most countable union of intervals if q is a stable fixed point of some f_n , while Ω_q is a finite or at most countable set if q is an unstable fixed point. The following lemma guarantees that flat parts of the graph of f_n grow up.

LEMMA 2.3.

- (i) *$\text{Fix}(f_n) \subset \text{Fix}(f_{n+1})$ for each n .*
 - (ii) *$f_n^{-1}(\text{Fix}(f_n)) \subset f_{n+1}^{-1}(\text{Fix}(f_{n+1}))$ for each n .*
- (Proof.) If $f_n(x) = x$, then $f_{n+1}(x) = (1 - \varepsilon)f_n(x) + \varepsilon f_n(f_n(x)) = (1 - \varepsilon)x + \varepsilon x = x$. Hence, $x \in \text{Fix}(f_{n+1})$.

Next, if $f_n(x) = q$ and $f_n(q) = q$ then $f_{n+1}(q) = q$ by (i) and $f_{n+1}(x) = (1 - \varepsilon)f_n(x) + \varepsilon f_n(f_n(x)) = (1 - \varepsilon)q + \varepsilon f_n(q) = q$.

Hence, $x \in f_{n+1}^{-1}(\text{Fix}(f_{n+1}))$. \square

(Proof of Theorem 2.1) Define the sets Ω and Ω^I by (2.3) and (2.4). If $x \in \Omega^I$, then $x \in \text{Fix}(f_n)$ for some n , and therefore, by lemma 2.3(i), $f_n(x) = x$ for all $m \geq n$. Hence, the limit $f_\infty(x)$ exists and equals x . If $x \in \Omega$, then $x \in f_n^{-1}(\text{Fix}(f_n(x)))$ for some n . By Lemma 2.3(ii) (and its proof)

$$f_m(x) = q \text{ for all } m \geq n \text{ with } q = f_n(x) \in \text{Fix}(f_n).$$

Hence, the limit $f_\infty(x)$ exists and equals q . In particular, $f_\infty(f_\infty(x)) = f_\infty(q) = q = f_\infty(x)$. Consequently we obtain (i) and (ii). Now (iii) follows if we set

$$\Omega_q = f_\infty^{-1}\{q\} = \{x \in \Omega | f_\infty(x) = q\}.$$

□

The case where f_∞ is a continuous function is very restrictive:

PROPOSITION 2.4. *If $f : I \rightarrow I$ is continuous and satisfies $f \circ f = f$ on I , then $\text{Fix}(f)$ is an interval and $I = f^{-1}\text{Fix}(f)$.*

(Proof.) The condition $f \circ f = f$ on I implies $I = f^{-1}\text{Fix}(f)$. Suppose the contrary. Then we would find two fixed points q_0 and q_1 and a non fixed point x in between q_0 and q_1 . Let $q_0 < x < q_1$. Since f is continuous, $f[q_0, q_1] \supset [q_0, q_1]$. Hence the intermediate value theorem implies that there exist a point y in (q_0, q_1) such that $f(y) = x$. Then we would have $f(f(y)) = f(x) \neq x$, which contradicts $f \circ f = f$. □

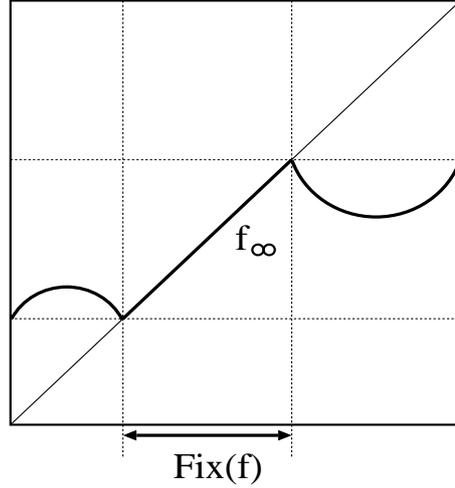


FIG. 2.1. *The case f_∞ is continuous.*

An example of such a f_∞ is shown in Figure 2.1. The continuous f_∞ is very restrictive, because even in a very simple case the function f_∞ is a non-continuous step function. The relevant f_∞ is actually a step function. Simple examples are as follows:

EXAMPLE 2.5. *If $f_0(x)$ is monotone nondecreasing and continuous, then $\Omega = I$ and $\text{Fix}^{(s)}(f_0) \subset \Omega^I \subset \text{Fix}(f_\infty) = \text{Fix}(f_0)$ (See Figure 2.2).*

(Proof.) Let J be a connected component of the set $\{x \in I \mid f_0(x) > x\}$. Then, $f_0(J) \subset J$ by the assumption on f_0 .

Thus, if $x \in J$ then the sequence $f_0(x), f_0(f_0(x)), \dots$ is monotone nondecreasing and bounded and so it has a limit $f_\infty(x)$, which is necessarily a fixed point of f_0 . In particular, the subsequence $\{f_n(x)\}_{n=0,1,2,\dots}$ converges to a fixed point of f_0 .

Similarly $f_n(x)$ converges to a fixed point of f_0 for any x is a connected components of $\{x \in I \mid f_0(x) < x\}$ □

EXAMPLE 2.6. *If the continuous map $f_0 : I \mapsto I$ has a stable fixed point whose basin of attraction is I or coincides with I except for an unstable fixed point, then $\Omega = I$ and $\Omega^I = \text{Fix}(f_\infty) = \text{Fix}(f_0)$.*

Now we consider the stable fixed points.

LEMMA 2.7. *If q is a (semi-)stable fixed point of f_0 , then q is a (semi-)stable fixed point of every f_n , $n \geq 0$.*

(Proof.) By definition, one can take a semi-open interval U which is either of the form $U = (q - \delta, q]$ or $[q, q + \delta)$ with $\delta > 0$, $U \supset f_0(U) \supset \dots \supset f_0^k(U) \rightarrow \{q\}$ as $k \rightarrow \infty$ and

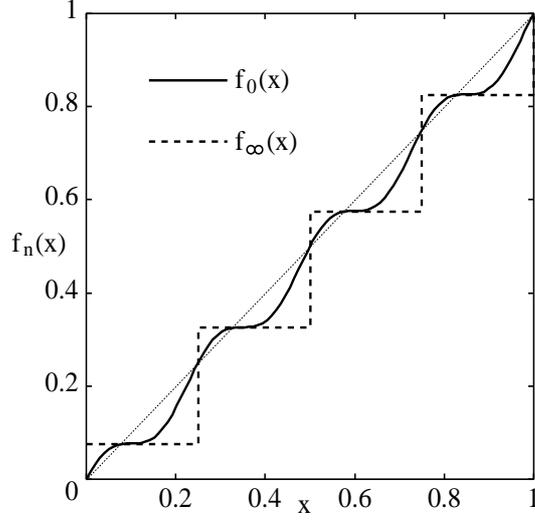


FIG. 2.2. The evolution of the function when f_0 is monotone nondecreasing and bounded. $f_0(x)$ (solid line) and $f_\infty(x)$ (dotted line) are plotted.

f_0 is one-to-one on U . Then,

$$\begin{aligned} f_1(U) &\subset (1 - \varepsilon)f_0(U) + \varepsilon f_0(f_0(U)) \\ &\subset (1 - \varepsilon)f_0(U) + \varepsilon f_0(U) \subset f_0(U) \subset U \end{aligned}$$

where $(1 - \varepsilon)A + \varepsilon B := \{(1 - \varepsilon)a + \varepsilon b \mid a \in A, b \in B\}$.

Similarly, for each $k \geq 1$, $f_1(f_0^k(U)) \subset f_0^{k+1}(U) \subset f_0^k(U)$.

Hence, we get

$$U \supset f_1(U) \supset \cdots \supset f_1^k(U)$$

and on the other hand

$$f_1^k(U) \subset f_0^k(U).$$

Consequently,

$$U \supset f_1(U) \supset \cdots \supset f_1^k(U) \rightarrow \{q\}.$$

The assertion for $n \geq 2$ follows by induction. \square

Now we discuss the degree of stability of fixed points. Let q be a fixed point of f_n , U be an interval contain q . Set

$$l_n(x; q) = l_n(x) = \frac{f_n(x) - q}{x - q} \quad \text{for } x \in U \text{ and } x \neq q.$$

Then,

$$(2.5) \quad l_{n+1}(x) = (1 - \varepsilon)l_n(x) + \varepsilon l_n(f_n(x))l_n(x).$$

Define $\rho_n = \rho_n(U)$ by $\rho_n = \sup_{x \in U, x \neq q} |l_n(x)|$, we get

$$(2.6) \quad 0 \leq \rho_{n+1} \leq (1 - \varepsilon)\rho_n + \varepsilon\rho_n^2.$$

This inequality shows that the dynamical system $h(x) = (1 - \varepsilon)x + \varepsilon x^2$ controls the stability of fixed points of f_n . From this observation we can show the following.

LEMMA 2.8. *Let q be a (semi-)stable fixed point of f_0 and U_q be a semi-open interval U_q which is either of the form $U_q = (q - \delta, q]$ or $[q, q + \delta)$ with $\delta > 0$, $U_q \supset f_0(U_q) \supset \dots \supset f_0^k(U_q) \rightarrow \{q\}$ as $k \rightarrow \infty$ and f_0 is one-to-one on U . If $\rho_0(U_q) < 1$, then $\lim_{n \rightarrow \infty} f_n(x) = q$.*

(Proof.) If $f_n(x) = q$ for some n , $\lim_{n \rightarrow \infty} f_n(x) = q$.

Suppose $f_n(x) \neq q$ for all n , then

$$\begin{aligned} |f_{n+1}(x) - q| &= |(1 - \varepsilon)(f_n(x) - q) + \varepsilon(f_n(f_n(x)) - q)| \\ &= \left| (1 - \varepsilon)(f_n(x) - q) + \varepsilon(f_n(x) - q) \frac{f_n(f_n(x)) - q}{f_n(x) - q} \right| \\ &= |f_n(x) - q| \cdot |(1 - \varepsilon) + \varepsilon l_n(f_n(x))| \end{aligned}$$

By inequality 2.6 and $\rho_0 < 1$,

$$|l_n(f_n(x))| < \rho_n < h(\rho_{n-1}) < h^n(\rho_0) \rightarrow 0 \text{ as } n \rightarrow \infty.$$

As a result, there exist δ such that $0 \leq |(1 - \varepsilon) + \varepsilon l_n(f_n(x))| < \delta < 1$ for large n and we have $|f_{n+1}(x) - q| < \delta |f_n(x) - q|$. This implies $\lim_{n \rightarrow \infty} f_n(x) = q$. \square

LEMMA 2.9. *For $x \in f_0^{-1}(U_q)$, $\lim_{n \rightarrow \infty} f_n(x) = q$.*

(Proof.) In the proof of lemma 2.8, a key inequality is $|l_n(f_n(x))| < \rho_n$ for $f_n(x) \in U_q$. Now $f_0(x) \in U_q$ from assumption. So $f_n(x) \in U_q$ by induction. The proof is similar. \square

From Lemmas 2.7-2.9, we can extend Theorem 2.1 as follows:

PROPOSITION 2.10. *For each n and $q \in \text{Fix}^{(s)}(f_n)$ take any interval V_q where f_n is monotone and such that $\rho_n(V_q) < 1$. Set*

$$\Omega = \bigcup_{n \geq 0} f_n^{-1} \left(\bigcup_{q \in \text{Fix}^{(s)}(f_n)} V_q \right).$$

Then the set Ω satisfies the conditions (i), (ii) and (iii) in Theorem 2.1.

REMARK 2.11. *There exist some unstable fixed points which become stable after iteration. For example, let q be a unstable fixed point of f_0 and assume that f_0 be monotone decreasing on U_q , $U_q \supset f_0(U_q)$ and $-1/\varepsilon < -\rho_0 < -(1 - \varepsilon)/\varepsilon$. Then q becomes a stable fixed point of f_1 because $\rho_1 < 1$ as shown in Figure 2.3.*

3. Trajectories driven by g_∞ and Hierarchical Map. Let us study the asymptotic behavior of trajectories other than (Ω, f_∞) . For this purpose, as is mentioned in Introduction, we write

$$(3.1) \quad f_{n+1} = g_n \circ f_n$$

by setting

$$(3.2) \quad g_n(x) = (1 - \varepsilon)x + \varepsilon f_n(x).$$

Since our target is the asymptotic behavior as $n \rightarrow \infty$, we may assume that $f_n|_\Omega$ is close to f_∞ from the beginning. For simplicity, we assume

$$(3.3) \quad f_0|_\Omega = f_\infty.$$

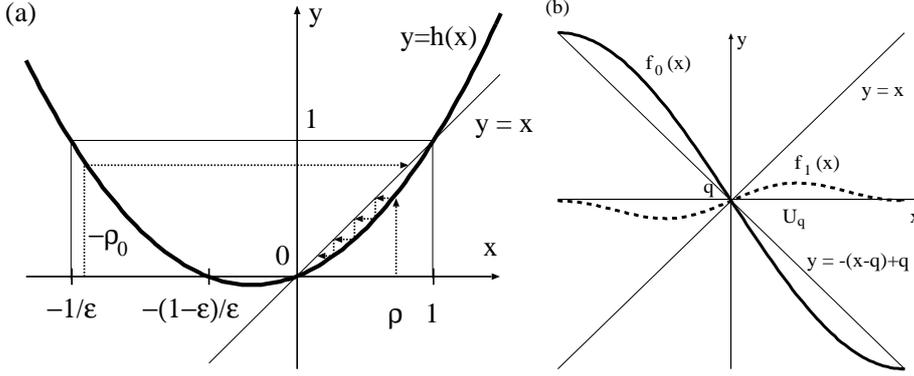


FIG. 2.3. (a) Dynamics of ρ_n and its example. (b) $f_0(x) = -\sin(\pi/2x)$ and $\varepsilon = 0.5$.

LEMMA 3.1. *If (3.3) holds, then*

$$(3.4) \quad f_n|_{\Omega} = f_{\infty}$$

for all $n = 0, 1, 2, \dots$

(Proof.) Assume $x \in \Omega$ and $f_n(x) = q \in \text{Fix}(f_{\infty})$. Then,

$$f_{n+1}(x) = (1 - \varepsilon)q + \varepsilon f_n(q) = q = f_n(x) = f_{\infty}(x).$$

Hence, $f_n|_{\Omega} = f_{\infty}$ implies $f_{n+1}|_{\Omega} = f_{\infty}$ and (3.4) follows by induction on n . \square

By Lemma 3.1, the “generated map” g_n also coincides with g_{∞} , if it is restricted to Ω :

$$(3.5) \quad g_{\infty}(x) = (1 - \varepsilon)x + \varepsilon q \text{ if } x \in \Omega_q.$$

EXAMPLE 3.2 (Nagumo-Sato map). *Let the initial map $f_0 : I \rightarrow I$, $I \supset [\frac{a-1}{\varepsilon}, \frac{a}{\varepsilon}]$ be as follows:*

$$f_0(x) = \begin{cases} \frac{a}{\varepsilon}, & x \in \{\frac{a}{\varepsilon}\} \cup (\frac{a-1}{\varepsilon}, 0), \\ \frac{a-1}{\varepsilon}, & x \in [0, \frac{a}{\varepsilon}) \cup \{\frac{a-1}{\varepsilon}\}, \\ b(x), & x \in I \setminus [\frac{a-1}{\varepsilon}, \frac{a}{\varepsilon}]. \end{cases}$$

Here $0 < a < 1$ and $b(x)$ is a function satisfying a condition $b(x) \in (\frac{a-1}{\varepsilon}, \frac{a}{\varepsilon})$. Then

$$\Omega = \left[\frac{a-1}{\varepsilon}, \frac{a}{\varepsilon} \right]$$

and

$$f_{\infty}(x) = \begin{cases} \frac{a}{\varepsilon}, & x \in \{\frac{a}{\varepsilon}\} \cup (\frac{a-1}{\varepsilon}, 0), \\ \frac{a-1}{\varepsilon}, & x \in [0, \frac{a}{\varepsilon}) \cup \{\frac{a-1}{\varepsilon}\}. \end{cases}$$

For instance suppose $f_0(x)$ is given by the Figure 3.1 (b). Then f_n converges to the f_{∞} shown in Figure 3.1(a). In this case, the generated map $g_{\infty}(x)$ is a piecewise linear map defined by:

$$g_{\infty} = \begin{cases} (1 - \varepsilon)x + a, & x \in \{\frac{a}{\varepsilon}\} \cup (\frac{a-1}{\varepsilon}, 0), \\ (1 - \varepsilon)x + a - 1, & x \in [0, \frac{a}{\varepsilon}) \cup \{\frac{a-1}{\varepsilon}\} \end{cases}$$

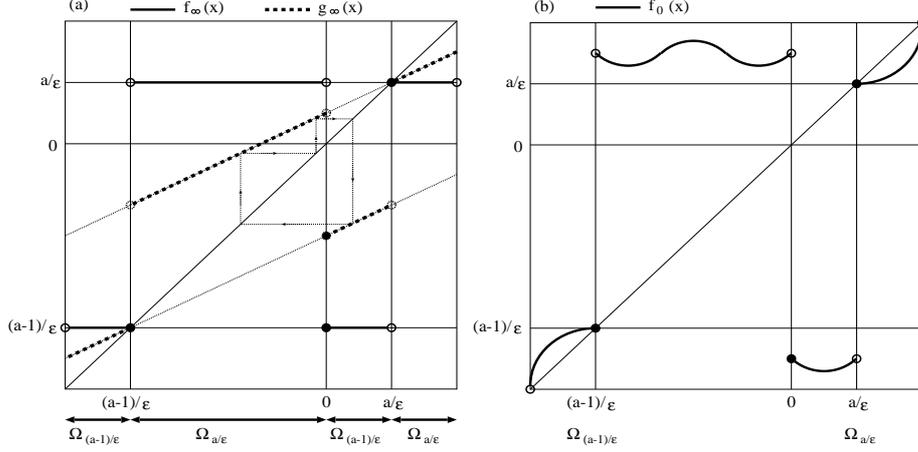


FIG. 3.1. (a) The graph of $f_\infty(x)$ (solid line) and the generated map $g_\infty(x)$ (dotted line). Here $\varepsilon = 1/2$ and $a = 1/4$. The map g_∞ has a period-3 attractor. (b) An example of a graph of initial function $f_0(x)$ which converges to f_∞ displayed in (a).

(See Figure 3.1(a)).

In this case, $g_\infty(\Omega) \subset \Omega$ and the map $g_\infty : \Omega \mapsto \Omega$ is studied in [3] and is called Nagumo-Sato map. It is shown that a periodic orbit of any period can appear if one controls ε and a suitably. Moreover, Cantor attractors (almost periodic orbits) can appear.

Combining (3.1), (3.2) and (3.5) we see $f_{n+1}(x) = g_\infty(f_n(x))$ if $f_n(x) \in \Omega$. Thus we obtain the trajectories $f_n(x)$ in Ω , which is driven by the Nagumo-Sato map g_∞ .

The above Example 3.2 is generalized to the following.

THEOREM 3.3. Assume $f_0|_\Omega = f_\infty$. Define $g_\infty : \Omega \mapsto I$ by (3.5). Set

$$(3.6) \quad \Omega(g_\infty) = \bigcap_{n \geq 0} g_\infty^{-n}(\Omega) = \{x \in \Omega \mid g_\infty(x), g_\infty(g_\infty(x)), \dots \in \Omega\}$$

and assume $\Omega(g_\infty)$ is non-empty. Then, the trajectories starting from $f_0^{-1}(\Omega(g_\infty))$ are “driven” by g_∞ . Precisely, if $x \in I$ and $f_0(x) \in \Omega(g_\infty)$ then

$$(3.7) \quad f_n(x) \in \Omega(g_\infty) \text{ for all } n = 0, 1, 2, \dots$$

and

$$(3.8) \quad f_n(x) = \underbrace{g_\infty \circ \dots \circ g_\infty}_{n \text{ times}}(f_0(x)).$$

In other words, the trajectory $\{f_n(x)\}_{n=0,1,2,\dots}$ on $\Omega(g_\infty)$ is reduced to the g_∞ -orbit of $f_0(x)$.

(Proof.) If we show (3.7), then (3.8) will be obvious from (3.5). Assume $f_n(x) \in \Omega(g_\infty)$, then $f_{n+1}(x) = g_\infty(f_n(x))$.

By the definition (3.6), $g_\infty(f_n(x)) \in g_\infty(\Omega(g_\infty)) \subset \Omega$.

Hence $f_{n+1}(x) \in \Omega(g_\infty)$. \square

Now we proceed to the next stage and seek for the set Ψ such that $f_n(\Psi) \subset \Omega$ and $g_n(\Psi) \subset \Psi$ for all n .

THEOREM 3.4. *Assume that $f_0|_\Omega = f_\infty$ and that there exist non empty f_∞ -invariant subsets Ω_i , $i = 1, \dots, N$, of $\Omega(g_\infty)$ and closed subsets Ψ_i , $i = 1, \dots, N$, of the complement of $\text{Fix}(f_\infty)$ which satisfies the following property: For each i there exists a j such that*

$$(3.9) \quad (1 - \varepsilon)\Psi_i + \varepsilon\Omega_i \subset \Psi_j.$$

Let $X_i = f_0^{-1}(\Psi_i)$. If $f_0(\Psi_i) \subset \Omega_i$ and $g_0(\Omega_i) \subset \Omega_i$, then $g_0(\Psi_i) \subset \Psi_j$. Moreover, for any $n \geq 1$,

$$(3.10) \quad f_n(X_i) \subset \Psi_i, f_n(\Psi_i) \subset \Omega_i \text{ and } g_n(\Psi_i) \subset \Psi_j.$$

Hence, once a trajectory $f_n(x)$ falls into some Ψ_i , say, $f_{n_0}(x) \in \Psi_i$, the trajectory $f_n(x)$, $n = n_0 + 1, n_0 + 2, \dots$ is confined in $\bigcup \Psi_j$ and driven by g_n 's: $f_{n+1} = g_n(f_n(x))$ for $n = n_0 + 1, n_0 + 2, \dots$

(Proof.) By (3.9) and the definition of g_0 ,

$$g_0(\Psi_i) \subset (1 - \varepsilon)\Psi_i + \varepsilon f_0(\Psi_i) \subset (1 - \varepsilon)\Psi_i + \varepsilon\Omega_i \subset \Psi_j.$$

Now, $f_1(X_i) = g_0(f_0(X_i)) \subset g_0(\Psi_i) \subset \Psi_j$ and $f_1(\Psi_i) = g_0(f_0(\Psi_i)) \subset g_0(\Omega_i)$. By Lemma 3.1, $g_0|_\Omega = g_\infty|_\Omega$. Thus $g_0(\Omega_i) \subset \Omega_i$. Hence (3.10) follows by induction. \square

REMARK 3.5. *If there exist closed subsets $\{X_i\}_{i=1}^M$ such that*

$$(3.11) \quad (1 - \varepsilon)X_i + \varepsilon\Psi_i \subset X_j \text{ for some } j,$$

the dynamics of $f_n|_{\Phi_i}$ ($f_n(\Phi_i) \subset X$) is determined by $g_n|_X$. This process can continue ad infinitum and it is not difficult to extend the Theorem 3.4.

In [2] the generated map $g_n|_\Psi$ is called the meta-map, taking into consideration the point that the dynamics of $g_n|_\Phi$ is determined by g_∞ . Similarly, in [2], the generated map $g_n|_{X_i}$ is called meta-meta-map, while the generated map g_n is called hierarchical map as a whole.

Now we present two typical examples.

The first example (Example 3.6) shows typical trajectories driven by g_∞ . On Ψ_i there are two branches of f_n , f_{even} and f_{odd} while on Ω_i f_∞ exists. The dynamics of $f_n|_{X_i}$ is determined by $g_n|_{\Psi_j}$.

The second example illustrates the case of f_n with further two branches on X_i . $f_n|_\Psi$ is driven by g_∞ , while $f_n|_X$ is driven by $g_n|_\Psi$, and $f_n|_\Phi$ is driven by $g_n|_X$ hierarchically. Each partial function is period 2 or time-invariant. The configuration of initial function f_0 is given by nesting the initial function of the first Example 3.6.

EXAMPLE 3.6 (meta-map). *In this example, a new initial function which generates a meta-map is shown. This initial function is given by a 'surgery' of the f_∞ which generates a map having a 2-period attractor. The partial function $f_n|_\Psi$ ($f_n(\Psi) \subset \Omega$) is set to generate a time-dependent $g_n|_\Psi$ which has another period-2 attractor (meta-map).*

Let $f_0(x)$ be as follows;

$$f_0(x) = \begin{cases} \begin{array}{ll} -(a+b), & x \in \Omega_{-(a+b)}, & \Omega_{-(a+b)} := \{-(a+b)\} \cup (-a, -(a-b)), \\ -(a-b), & x \in \Omega_{-(a-b)}, & \Omega_{-(a-b)} := \{-(a-b)\} \cup (-(a+b), -a], \\ a-b, & x \in \Omega_{a-b}, & \Omega_{a-b} := \{a-b\} \cup (a, a+b), \\ a+b, & x \in \Omega_{a+b}, & \Omega_{a+b} := \{a+b\} \cup (a-b, a], \end{array} \\ \\ \begin{array}{ll} -a+Eb, & x \in \Psi_0, & \Psi_0 := (0, a-b), \\ a+Eb, & x \in \Psi_1, & \Psi_1 := (-(a-b), 0], \end{array} \\ \\ \begin{array}{ll} -E(a+Eb), & x \in X_0, \\ E(a-Eb), & x \in X_1. \end{array} \end{cases}$$

Here $E := \frac{\varepsilon}{2-\varepsilon}$ and $b = \frac{1-\varepsilon}{1+\varepsilon} \cdot a$ ($a > 0$). The graph of this initial function is shown in Figure 3.2. The initial function f_0 on Ω_i is similar to the f_0 with two fixed points in Example 3.1, two copies of which are now embedded in subintervals $[-(a+b), -(a-b)]$ and $[a-b, a+b]$ for the initial function f_0 . Here, the function which generates a map g_∞ having period-2 attractor is embedded to the subintervals. Now, $\text{Fix}(f_0) = \{\pm(a+b), \pm(a-b)\}$.

Now, (i) The generated map g_∞ has a period-2 attractor. (ii) $f_0|_\Psi$ ($f_n(\Psi) \subset \Omega$ for all n) is on the attractor of g_∞ . (iii) The $f_0|_\Psi$ is arranged so as to generate a time-dependent map $g_n|_\Psi$ (meta-map), which has another period-2 attractor. (iv) $f_0|_X$ ($f_n(X) \subset \Psi$ for all n) is on the attractor of $g|_\Psi$. Each partial function is already on one of the attractors and f_n is a period-2 function as a whole.

The procedure of time evolution is demonstrated straightforwardly as follows (For the computation of each step, it is convenient to use the relation $(1-\varepsilon)Ea + \varepsilon(-a) = -Ea$).

At $n = 0$, the following conditions are satisfied.

$$\begin{cases} f_0(\Psi_0) \subset \Omega_{-(a+b)}, \\ f_0(\Psi_1) \subset \Omega_{a-b}, \\ \\ f_0(X_0) \subset \Psi_1, \\ f_0(X_1) \subset \Psi_0. \end{cases}$$

At the next step, this $f_0(x)$ evolves to the following f_1 ;

$$\begin{cases} \begin{array}{ll} f_1|_{\Psi_0} = (1-\varepsilon)(-a+Eb) + \varepsilon(-(a+b)) & = & -a - Eb, \\ f_1|_{\Psi_1} = (1-\varepsilon)(a+Eb) + \varepsilon(a-b) & = & a - Eb, \end{array} \\ \\ \begin{array}{ll} f_1|_{X_0} = (1-\varepsilon)(-E(a+Eb)) + \varepsilon(a+Eb) & = & E(a+Eb), \\ f_1|_{X_1} = (1-\varepsilon)(E(a-Eb)) + \varepsilon(-a+Eb) & = & -E(a-Eb). \end{array} \end{cases}$$

Then, at the step $n = 1$, the following conditions are satisfied. Note that there is an exchange of suffices of Ψ_i .

$$\begin{cases} \begin{array}{ll} f_1(\Psi_0) \subset \Omega_{-(a-b)}, \\ f_1(\Psi_1) \subset \Omega_{a+b}, \end{array} \\ \\ \begin{array}{ll} f_1(X_0) \subset \Psi_0, \\ f_1(X_1) \subset \Psi_1. \end{array} \end{cases}$$

At the next step, the f_1 evolves to the following f_2 ;

$$\begin{cases} f_2|_{\Psi_0} = (1-\varepsilon)(-a-Eb) + \varepsilon(-a-b) = -a+Eb = f_0|_{\Psi_0}, \\ f_2|_{\Psi_1} = (1-\varepsilon)(a-Eb) + \varepsilon(a+b) = a+Eb = f_0|_{\Psi_1}, \\ f_2|_{X_0} = (1-\varepsilon)E(a+Eb) + \varepsilon(-a-Eb) = -E(a+Eb) = f_0|_{X_0}, \\ f_2|_{X_1} = (1-\varepsilon)(-E(a-Eb)) + \varepsilon(a-Eb) = E(a-Eb) = f_0|_{X_1}. \end{cases}$$

This f_2 coincides with f_0 . Hence f_n is a period-2 function. These dynamics are shown in Figure 3.2, while Figure 3.3 shows a schematic representation of the dynamics. Each arrow $A \rightarrow B$ in the figure indicates that $f_n(A) \subset B$. $f_\infty(\Omega_i)$ is always included in $\text{Fix}(f_\infty)$ and $f_n(\Psi_0)$ is included in Ω_{-1} or $\Omega_{-\varepsilon}$ in turns.

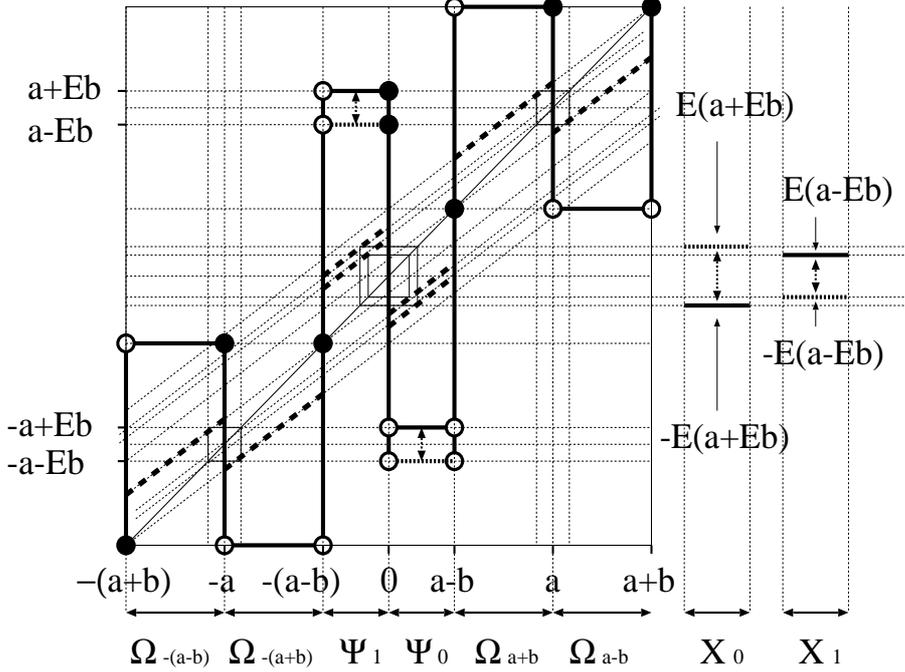


FIG. 3.2. Temporal evolution of f_n with a meta-map. Here, $E := \frac{\varepsilon}{2-\varepsilon}$ and $b = \frac{1-\varepsilon}{1+\varepsilon} \cdot a$. $f_n(x)$ at even step n is plotted by solid line, while that for odd n is plotted by dotted line. $g_n(x)$ is shown by bold dotted line. ε is set to $1/4$.

In the previous example, the function which generates a map having a period-2 attractor is embedded to give a new initial function f_0 . Note that the ‘surgery’ of the initial function is valid so that the generated map of the function has an arbitrary period.

The next example shows a meta-meta-map given by nesting this initial function.

EXAMPLE 3.7 (meta-meta-map). Define a new initial function f_0 by a recursive ‘surgery’ of the f_0 in Example 3.6 (meta-map). The meta-meta-map is given by this recursive surgery. In Figure 3.4, the hierarchical configuration of f_0 is plotted. Two copies of the initial function in the Example 3.6 (meta-map) are embedded on the intervals $[-(a+b), -(a-b)]$ and $[a-b, a+b]$, for this new f_0 . Now, f_0 has 8 fixed points $\text{Fix}(f_\infty) = \{\pm(a+b), \pm(a-b), a \pm c, -a \pm c\}$. Here, $c := \frac{a-b}{a+b} \cdot b$. According to the previous example, (i) each $g_n|_{\Phi}$ (meta-map) has period-2 attractors and (ii) $f_n|_X$

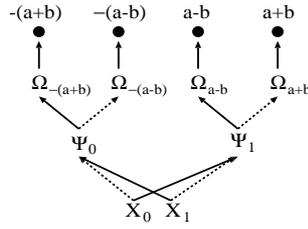


FIG. 3.3. The schematic representation of the dynamics. The arrow $A \rightarrow B$ means that $f_n(A) \subset B$. f_{even} (solid line) and f_{odd} (dotted line) are shown. The arrows over time steps (i.e., over the periods ($=2$)) are overlaid.

is arranged on the attractor of $g_n|_{\Psi}$. (iii) In this example, $f_n|_X$ is set to generate a time-dependent $g_n|_X$ which has period-2 attractors (meta-meta map). (iv) Each $f_n|_{\Phi_i}$ ($f_n|_{\Phi_i} \subset X$ for all n) is on the attractor of $g_n|_X$ and gives a period-2 function.

Figure 3.5 shows a schematic representation of this case. (i) $f_n|_{\Psi}$ is driven by $g|_{\Omega}$, (ii) $f_n|_X$ driven by $g_n|_{\Psi}$ and (iii) $f_n|_{\Phi}$ is driven by $g_n|_X$ hierarchically. As is shown in this figure, one more step ($f_n|_{\Phi}$) is added to the hierarchy in the Example 3.6 (Figure 3.3), here.

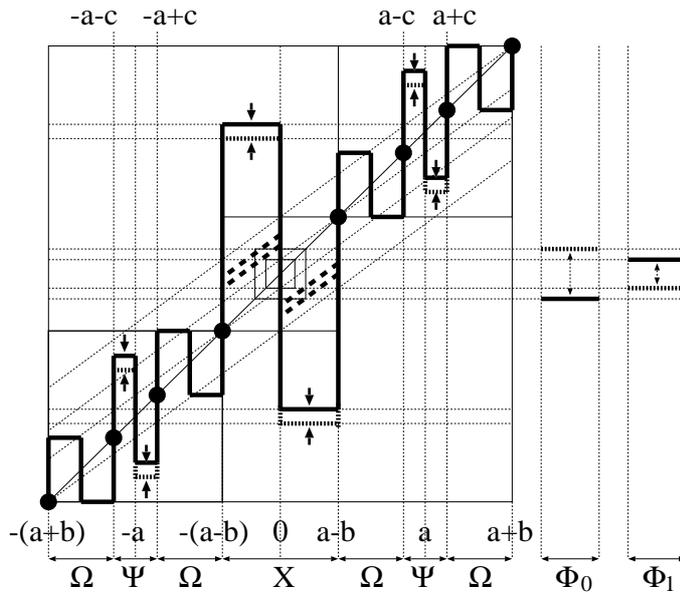


FIG. 3.4. The evolution of f_n . For the initial function, two copies of the Example 3.6 (Figure 3.2) are embedded to $[\pm(a+b), \pm(a-b)]$. Here, ε is set to $1/4$ and $c := \frac{a-b}{a+b} \cdot b$.

As is described in the remark 3.5, this process can be continued ad infinitum. A simple method to give an initial function with a higher hierarchical structure is to nest a given f_0 so that it satisfies the condition for the (extended) theorem 3.4.

4. Further Example. In the Examples 3.6 (meta-map) and 3.7 (meta-meta-map), the intervals are partially ordered at each step, if the order is defined so that $I_a < I_b$ iff $I_a \subset f_n(I_b)$ are satisfied (See Figure 3.3 and 3.5). In the Figures 3.3 and

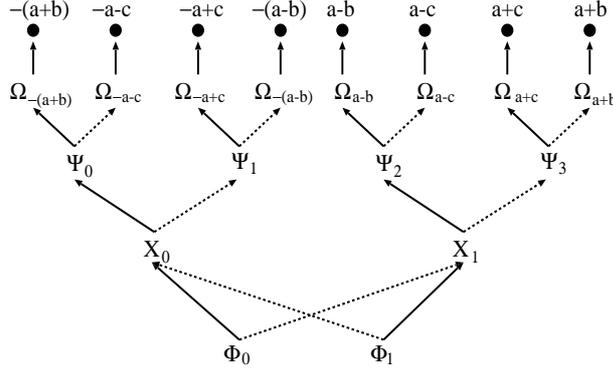


FIG. 3.5. The schematic representation of the dynamics. One more step ($f_n|_{\Phi}$) is added to the hierarchy in the Example 3.6 (Figure 3.3).

3.5, the arrows $A \rightarrow B$ for $f_n(A) \subset B$ change in time and the arrows over time steps (i.e., over the periods ($=2$)) are overlaid. Note that the intervals there are partially ordered. Generally, the intervals are not partially ordered for overlaid graph over n . An example of the initial function for such case is given below. In this example, the hierarchy is “entangled”. There are some partial functions driven by each other generated map in turns.

EXAMPLE 4.1 (entangled hierarchy). Let $f_0(x)$ be as follows;

$$f_0(x) = \begin{cases} \varepsilon - 3, & x \in \Omega_{\varepsilon-3}, & \Omega_{\varepsilon-3} := [\varepsilon - 3, -2] \cup (-1, 0], \\ 3 - \varepsilon, & x \in \Omega_{3-\varepsilon}, & \Omega_{3-\varepsilon} := (0, 1] \cup (2, 3 - \varepsilon], \\ 1 + \varepsilon, & x \in \Psi_0, & \Psi_0 := (-2, -1], \\ -(1 - \varepsilon), & x \in \Psi_1, & \Psi_1 := (1, 2]. \end{cases}$$

In this example, the time evolution of this f_0 is demonstrated directly as follows. Now, the initial function satisfies the following condition.

$$\begin{cases} f_0(\Psi_0) \subset \Psi_1, \\ f_0(\Psi_1) \subset \Omega_{\varepsilon-3}. \end{cases}$$

This f_0 evolves to the following f_1 .

$$\begin{cases} f_1|_{\Psi_0} = (1 - \varepsilon)f_0|_{\Psi_0} + \varepsilon f_0|_{\Psi_1} \circ f_0|_{\Psi_0} & = 1 - \varepsilon, \\ f_1|_{\Psi_1} = (1 - \varepsilon)f_0|_{\Psi_1} + \varepsilon f_0|_{\Omega_{\varepsilon-3}} \circ f_0|_{\Psi_1} & = -(1 + \varepsilon). \end{cases}$$

At $n = 1$ the relation

$$\begin{cases} f_1(\Psi_0) \subset \Omega_{3-\varepsilon}, \\ f_1(\Psi_1) \subset \Psi_0 \end{cases}$$

is satisfied. In this case, Ψ is not only mapped to Ω , but also to Ψ itself. This f_1 evolves to the following f_2

$$\begin{cases} f_2|_{\Psi_0} = (1 - \varepsilon)f_1|_{\Psi_0} + \varepsilon f_1|_{\Omega_{3-\varepsilon}} \circ f_1|_{\Psi_0} & = f_0|_{\Psi_0}, \\ f_2|_{\Psi_1} = (1 - \varepsilon)f_1|_{\Psi_1} + \varepsilon f_1|_{\Psi_0} \circ f_1|_{\Psi_1} & = f_0|_{\Psi_1}. \end{cases}$$

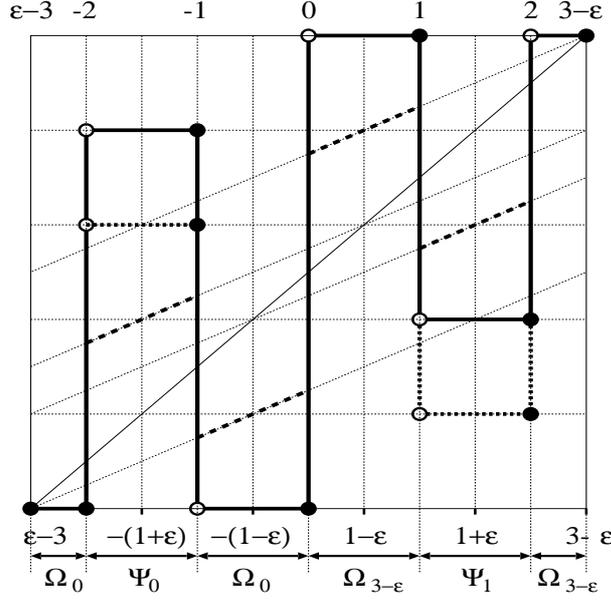


FIG. 4.1. The evolution of the function with entangled hierarchy. The solid line shows $f_{\text{even}}(x)$ and dotted line shows $f_{\text{odd}}(x)$. The bold dotted line shows $g_n(x)$. $f_n|_{\Psi_i}$ evolves with period 2. Here, $\varepsilon = 1/2$.

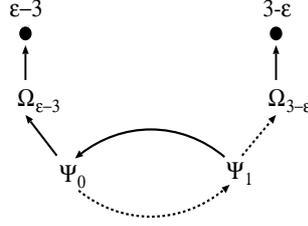


FIG. 4.2. The schematic representation of the “entangled hierarchy”. The arrows over time steps (i.e., over the periods (=2)) are overlaid. There is a loop in this graph which does not exist in previous examples.

This f_2 coincides with f_0 . Hence f_n is a period-2 function. These are shown in Figure 4.1 while the schematic representation is shown in Figure 4.2.

The loop in Figure 4.2 shows that the dynamics of $f_n|_{\Psi_0}$ and $f_n|_{\Psi_1}$ are determined by $g_n|_{\Psi_1}$ and $g_n|_{\Psi_0}$ in turns. Note that the dynamics of $f_n|_{\Psi_0}$ and $f_n|_{\Psi_1}$ are ‘not’ determined each other at the same step n in the Example 4.1. The snapshot of the graph at n is partially ordered, while the overlaid graph for n has the loop.

In Section 3, all intervals are partially ordered. There, the dynamics of $f_n|_A$ have no influence to the dynamics of $f_n|_B$, if $f_n(A) \subset B$. Now the “entanglement” exists and the dynamics of $f_n|_A$ has the influence to $f_m|_B$ ($n \neq m$), even if the condition $f_n(A) \subset B$ are satisfied at some n .

5. Discussion. To close the paper, we briefly discuss the original motivation in the study of (1.1) [1] and possible relevance of our result to a biological system. In a biological system, we are often amazed at its ability to change its own rule, while

in a mechanical system there usually exists a rigid rule which governs the change of the state forever. Moreover, the rule in a biological system is formed ‘spontaneously’, depending on the history of the state, without being prescribed externally. There a rule to drive the change of the state and the state driven by the rule are not separated initially, but through dynamics, some part of the system starts to drive other parts, and works as a rule.

When we adopt usual dynamical systems on phase spaces, however, the question how a rule is formed is not answered, since in dynamical systems, the rule for dynamics, and the variables that are driven by the rule are clearly separated. When a rule is not separated from the state, however, the rule (that is undifferentiated from the state variables) may operate to itself. In our function dynamics, we try to answer the problem of this self-operation of a rule by explicitly taking into account the term $f \circ f$, since with this term, the function f to change a state value x can also be a state value to be changed by it.

This $f \circ f$ term leads to a self-reference, since the evolution of the function $f_n(x)$ obeys the generated map $g_n(x)$, which itself refers to the function $f_{n-1}(x)$. Indeed the importance of self-reference is generally discussed in a biological problem. In our cognition, for example, external inputs are processed and are mapped to an output. The output from this process influences our cognitive process itself. If we regard this cognitive process as a function from inputs to outputs, this function changes in time following some self-reference, through development of our cognition. Our study of the function dynamics (1.1) was originally introduced as a toy model to study the dynamics with such self-reference [1], and was motivated by the search for a novel class of phenomena in a system with self-referential structure.

In the structure of Section 3, we have demonstrated that evolution of some partial functions is driven by the generated map of some other intervals hierarchically. The generated map of some intervals works as a ‘rule’ to drive other intervals, although they are not initially prescribed as a rule part in our model equation. These intervals to drive other parts are given by flat parts of $f_n(x)$. In fact, with temporal evolution of our function dynamics, the whole interval is partitioned into flat parts.

In a biological system, rules are often formed first by partition of continuous inputs into discrete symbols, and these symbols provide a basis for a syntactic structure to drive other parts. This partition process is called articulation in our cognition and language (for example, continuous spectrum of light is ‘articulated’ into a discrete set of colors). As mentioned, this articulation process and the generation of rules over the articulated symbols are a general feature of our function dynamics.

In the function dynamics, the rule, i.e., the generated map, can change in time, when the driving by a generated map has a hierarchical structure as in Section 3. In this sense, the hierarchy of a rule, a rule to change the rule, the further rule to change it, ... is formed in the function dynamics. Such hierarchy in the change of rules also reminds us of hierarchical structures ubiquitous in a biological system, and also in our cognitive process. Furthermore, in Section 4, we have found an example in which the hierarchy structure itself can change in time, where the separation of rule and state formed through dynamics is partially destroyed. The rules and hierarchy in a biological system have stability on one hand, and plasticity on the other hand. In future, it will be important to analyze the stability of the structure we found in the paper.

The late Professor Masaya Yamaguti stressed the importance of self-reference

from early days. He often mentioned his interest in fractals in connection with the self-reference. Hata and he studied function equation leading to fractal [4]. Furthermore he often discussed that mathematics dealing with self-reference is necessary to psychology, natural language, and so forth. It is to be regretted that we could not present our paper while he lived.

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